

# #1

ENGINYERIA  
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# APUNTS

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TÍTOL / **TRANSMISSION LINES AND RADIOFREQUENCY CIRCUITS ; ANTENNAS**

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CENTRE / EETAC

DATA / 2019

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# 1. TRANSMISSION LINES AND RADIOFREQUENCY CIRCUITS

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- Introduction
- Transmission lines
  - Transmission line types
  - Propagation modes
  - Propagation equations
    - lossless lines
    - lossy lines
    - low-loss lines
  - Reflection coefficient → forward and reverse propagating waves
  - Power and losses: return loss → attenuation
  - Voltage standing wave ratio
  - Impedance → impedance of transmission lines
  - Generator mismatch

[COLLIN] R.E. Collin, *Foundations for Microwave Engineering*, Wiley-Interscience, 2<sup>nd</sup> Edition, 2001 (New York)

[POZAR] D.M. Pozar, *Microwave Engineering*, Addison-Wesley Publishing Company, 2<sup>nd</sup> Edition, 1993 (Reading, Massachusetts)



## GLOSSARY

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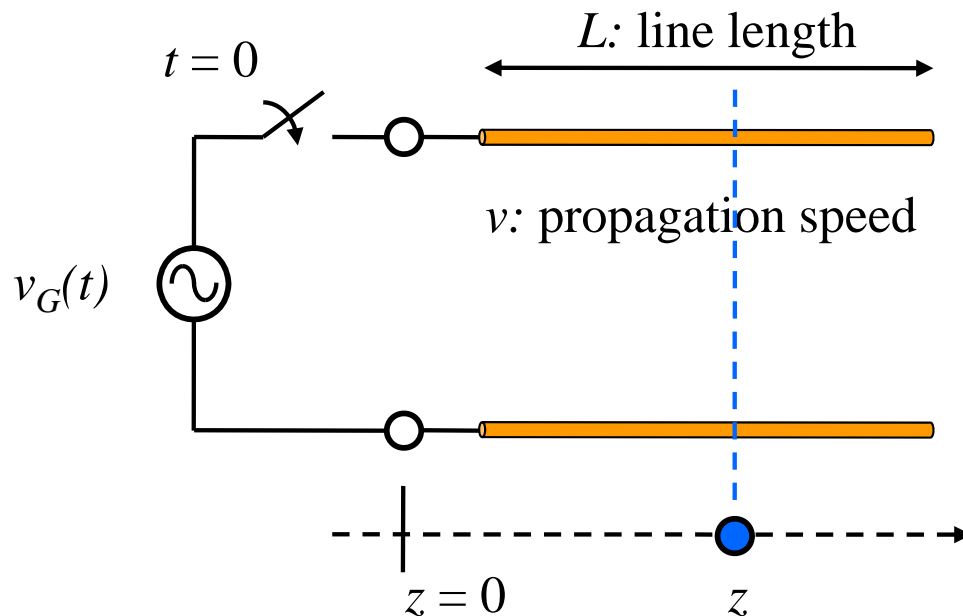
- $\alpha$  : attenuation constant [ $\text{m}^{-1}$ ]
- $\beta$  : phase constant [ $\text{rad}\cdot\text{m}^{-1}$ ]
- $C_d$  : distributed capacitance per unit length [ $\text{F}/\text{m}$ ]
- **CNS** : Communications, Navigation, Surveillance
- $\epsilon_0$  : electric permittivity of vacuum [ $8.85\cdot 10^{-12} \text{ F}/\text{m}$ ]
- $f_0$  : frequency [ $\text{Hz}$ ]
- $\gamma$  : propagation constant [ $\text{m}^{-1}$ ]
- $G$  : distributed conductance per unit length [ $\text{S}/\text{m}$ ]
- $i(z,t)$  : current in time domain [ $\text{V}$ ]
- $I_0^+$  : current amplitude of progressive wave at  $z=0$  [ $\text{A}$ ]
- $l$  : transmission line length [ $\text{m}$ ]
- $L_d$  : distributed inductance per unit length [ $\text{H}/\text{m}$ ]
- $\lambda$  : wavelength [ $\text{m}$ ]
- $\mu_0$  : magnetic permeability of vacuum [ $4\pi\cdot 10^{-7} \text{ H}/\text{m}$ ]
- $\omega$  : angular frequency [ $\text{rad}/\text{s}$ ]
- $R$  : distributed resistance per unit length [ $\Omega/\text{m}$ ]
- **RF** : radiofrequency
- $T$  : period [ $\text{s}$ ]

## GLOSSARY

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- $R_S$  : surface resistivity [ $\Omega$ /square]
- $\delta_S$  : skin depth [m]
- $\sigma$  : conductivity [Sm]
- $\sigma_d$  : dielectric conductivity [Sm]
- $\tan \delta$  : loss tangent [adim]
- $RL$  : return loss [dB]
- $\rho$  : (voltage) reflection coefficient [adim]
- $\rho_G$  : (voltage) generator reflection coefficient [adim]
- $\rho_{IN}$  : (voltage) reflection coefficient at input port [adim]
- $\rho_L$  : (voltage) load reflection coefficient [adim]
- $v(z,t)$  : voltage in time domain [V]
- $V_G$  : voltage at generator [V]
- $V_0^+$  : voltage amplitude of progressive wave at  $z=0$  [V]
- $v_p$  : phase velocity [m/s]
- $VSWR$  : Voltage Standing Wave Ratio [adim]
- $Z_G$  : generator impedance [ $\Omega$ ]
- $Z_{IN}$  : impedance at the input port of the transmission line [ $\Omega$ ]
- $Z_L$  : load impedance [ $\Omega$ ]
- $Z_0$  : transmission line characteristic impedance [ $\Omega$ ]

- Wave propagation along a transmission line considering the propagation delays:



Source voltage:

$$v_G(t) = A \cdot \cos(\omega_0 t)$$

Voltage:

$$v(z, t) = A \cdot \cos\left(\omega_0 \left[t - \frac{z}{v}\right]\right)$$

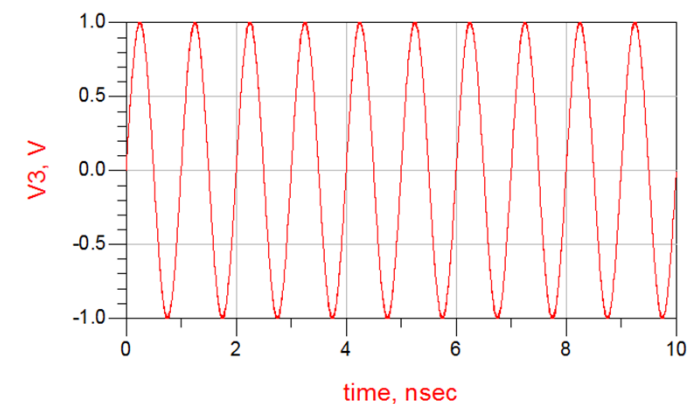
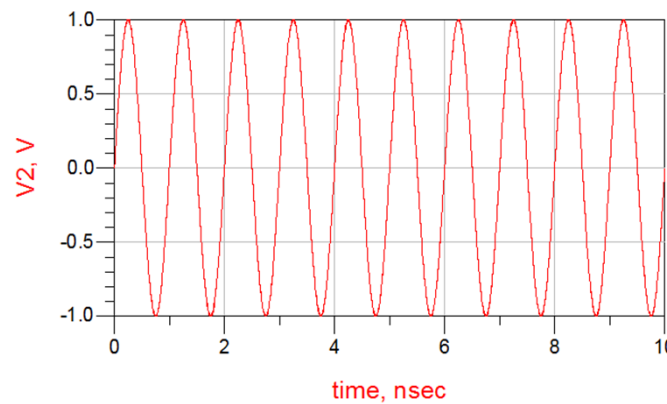
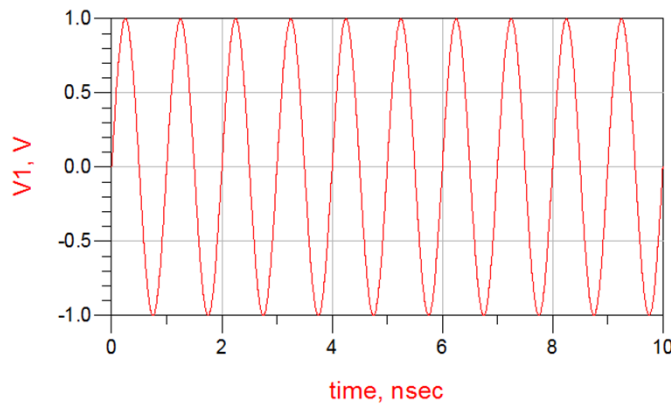
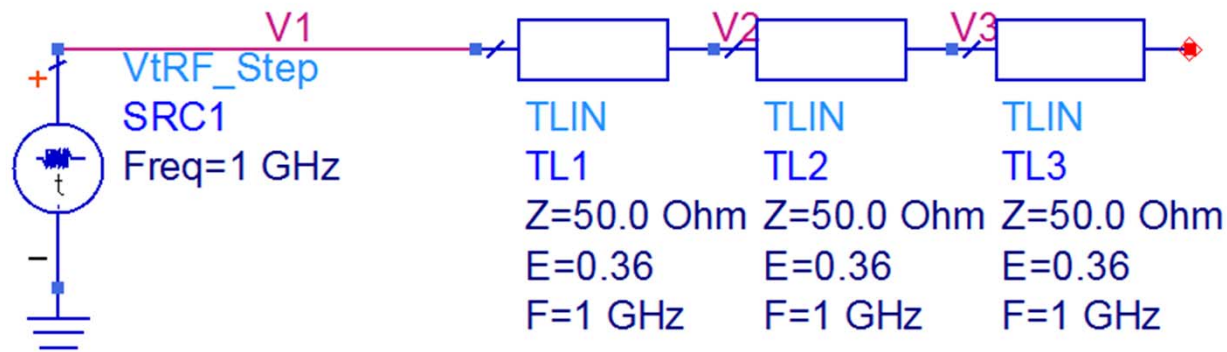
propagation delay

- Some remarks:

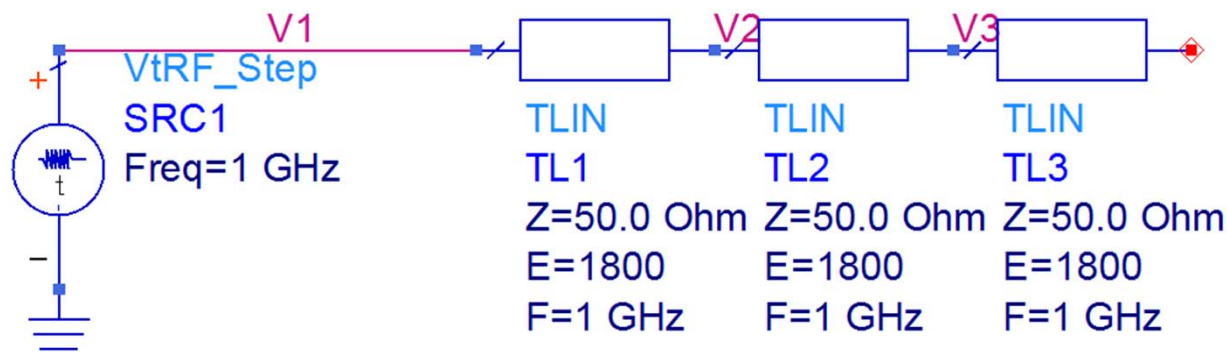
- Wave propagation considers the propagation delays.
- Each point in the line has a different voltage/current at the same time  $t$ .
- Periodicity in time or *period*  $T$ .
- Spatial periodicity or *wavelength*  $\lambda$ .
- Dimensions use to be define with respect to  $\lambda$ .

$$T = \frac{2\pi}{\omega_0} \quad \lambda = \frac{2\pi}{\omega_0/v} = \frac{v}{f_0}$$

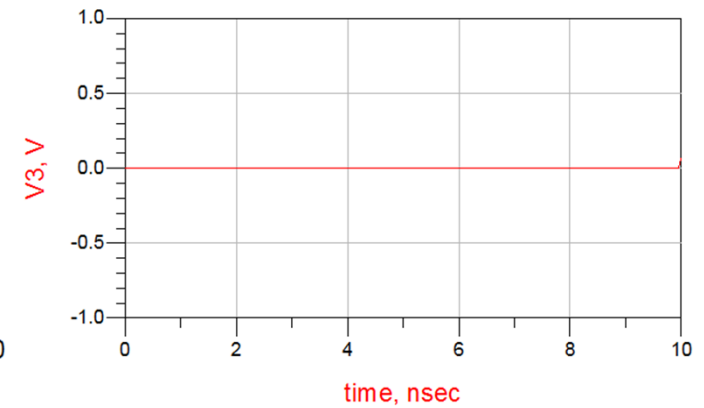
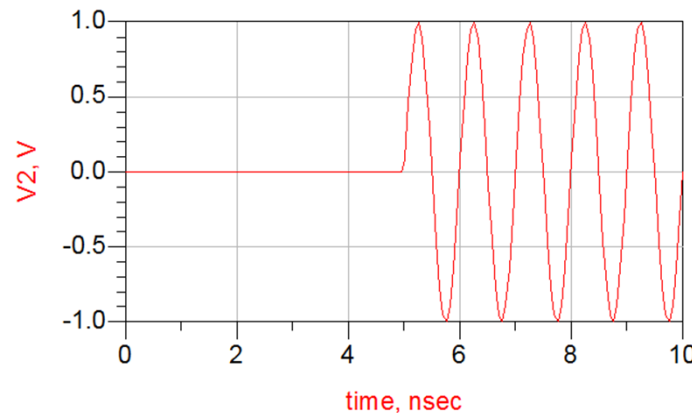
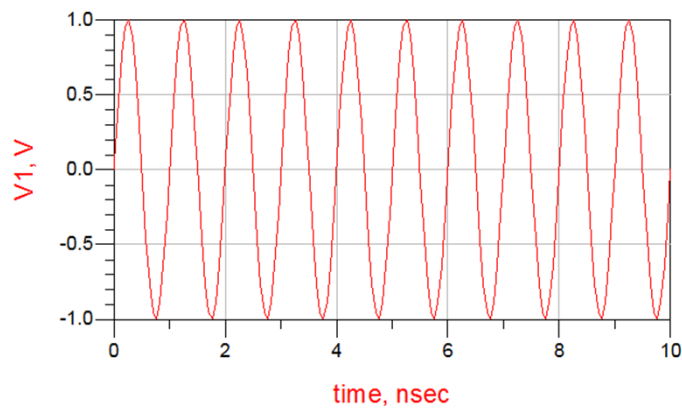
**Example: Propagation at Low Frequencies.** Consider a circuit having a transmission line length of  $0.003\lambda$  (3 lines,  $0.001\lambda$ -length each) fed with a sinusoidal wave of 1 GHz.



**Example: Propagation at High Frequencies.** Consider a circuit having a transmission line length of  $15\lambda$  (3 lines,  $5\lambda$ -length each) fed with a sinusoidal wave of 1 GHz.



Tran  
Tran1  
StopTime=10.0 nsec  
MaxTimeStep=0.1 nsec



# TRANSMISSION LINES

## Transmission line types

- Transmission line: a set of two or more than two conductors in a dielectric medium having a uniform transversal section that allow the propagation of a EM wave.
- Transmission lines are physical devices whose purpose is to guide electromagnetic waves (carry RF power) from one place to another.
- They are capable of guiding TEM waves (TEM waves can only exist in structures containing two or more separated conductors).
- Two-wire transmission lines are inefficient for transferring electromagnetic energy at high frequencies due to the lack of confinement in all directions.
- Coaxials are more efficient than two-wire lines in those cases.

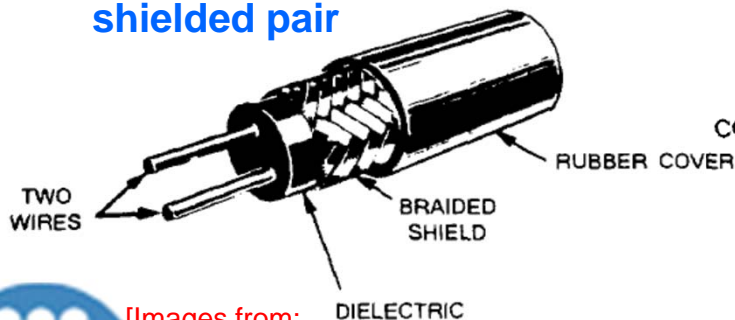
**two-wire ribbon line  
(twin lead)**



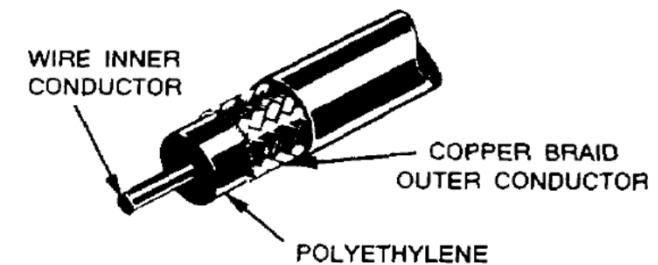
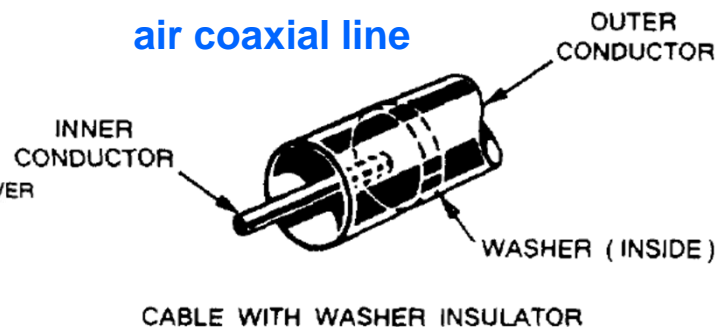
**twisted pair**



**shielded pair**



**air coaxial line**



**flexible coaxial line**

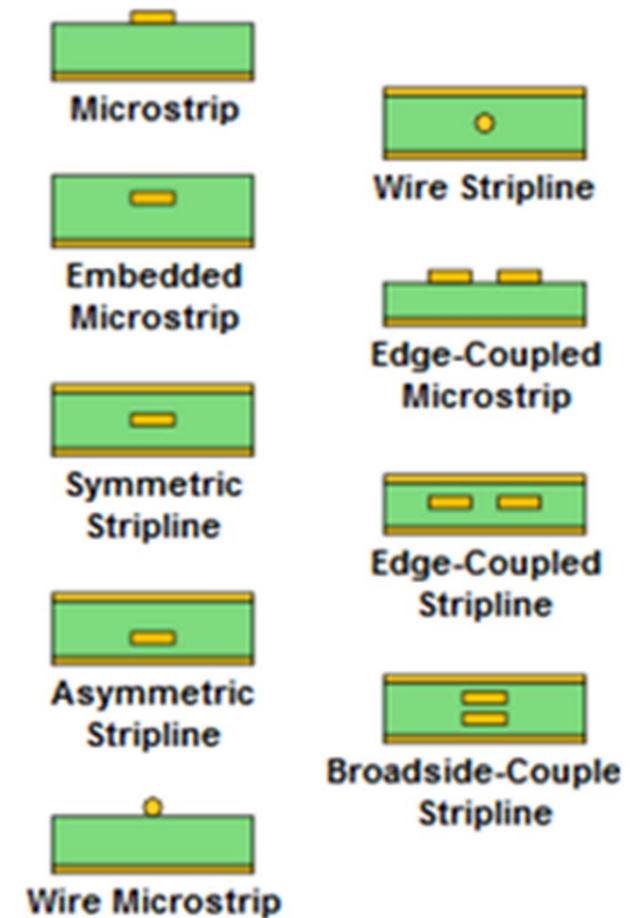
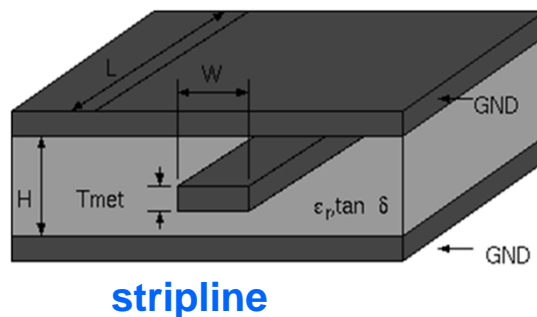
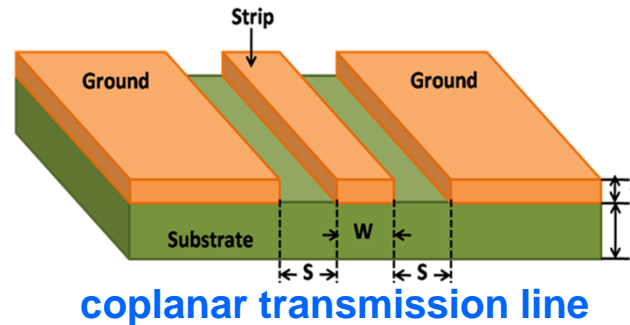
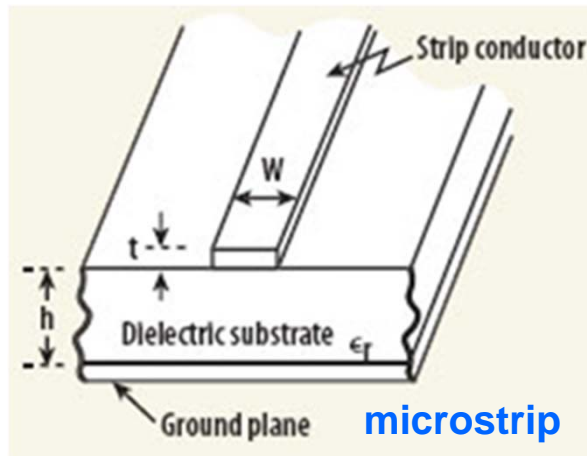
[Images from: <http://www.techlearner.com/Apps/TransandGuides.pdf>]



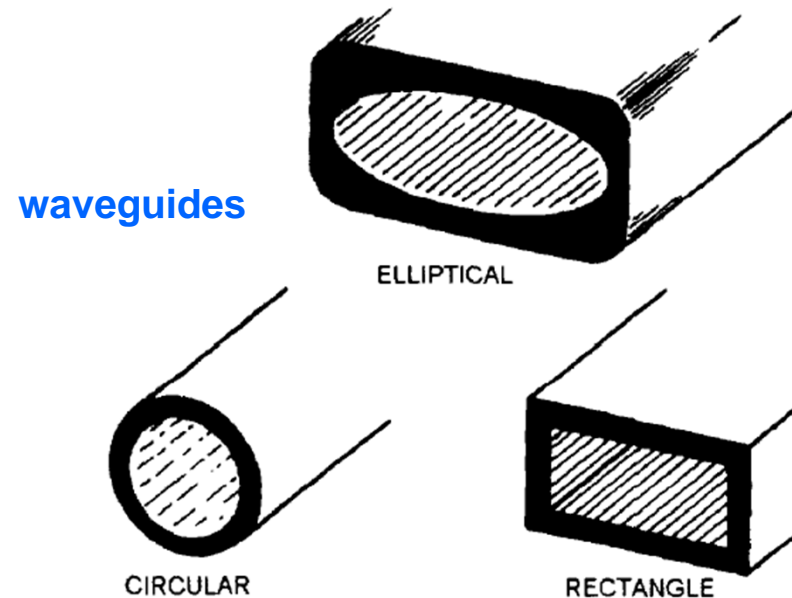
# TRANSMISSION LINES

## Transmission line types

- There are lots of planar structures used as transmission lines. Metallic parts are supported by dielectrics (fiberglass, ceramics, foams,...).

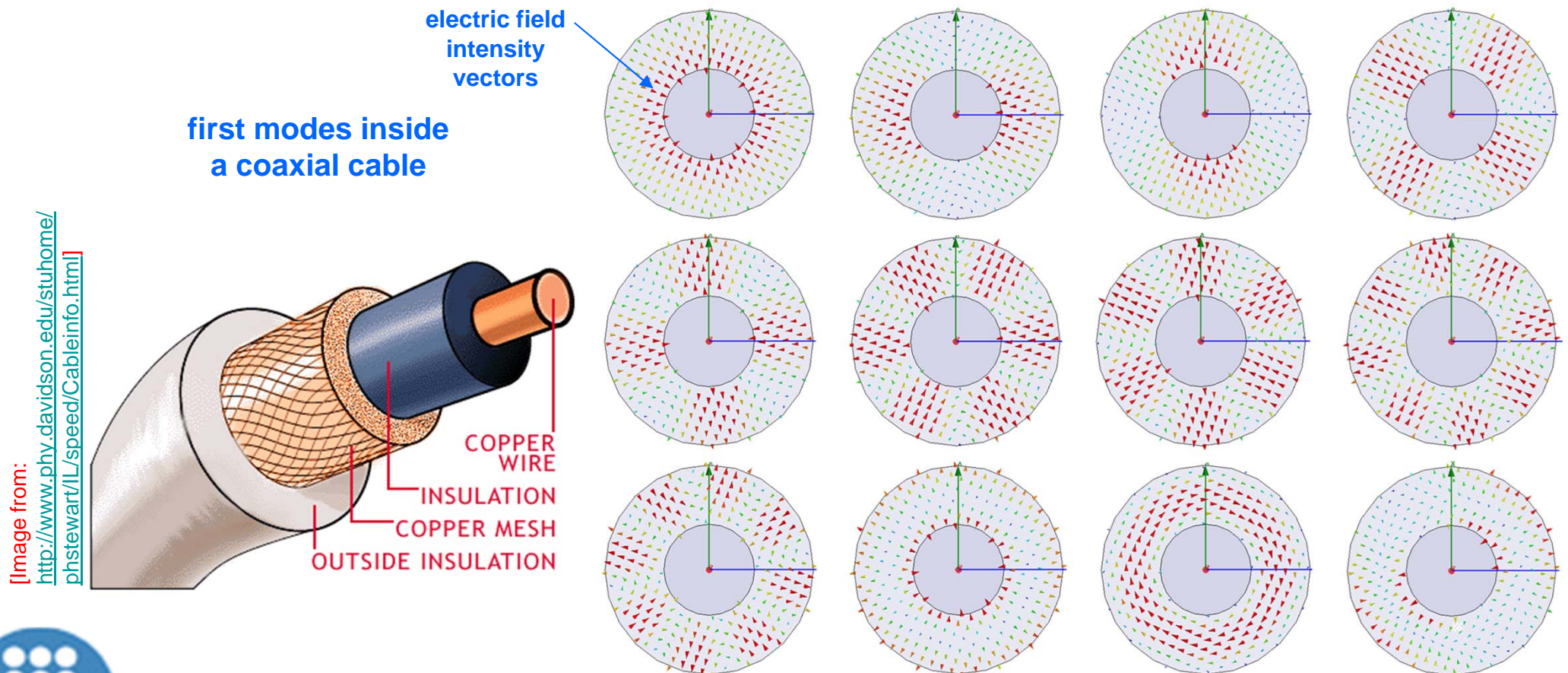


- Waveguides are the most efficient. They are fabricated with just one conductor. Waveguides do not support TEM waves.
- Two-wire lines are less bulky and less expensive than waveguides.



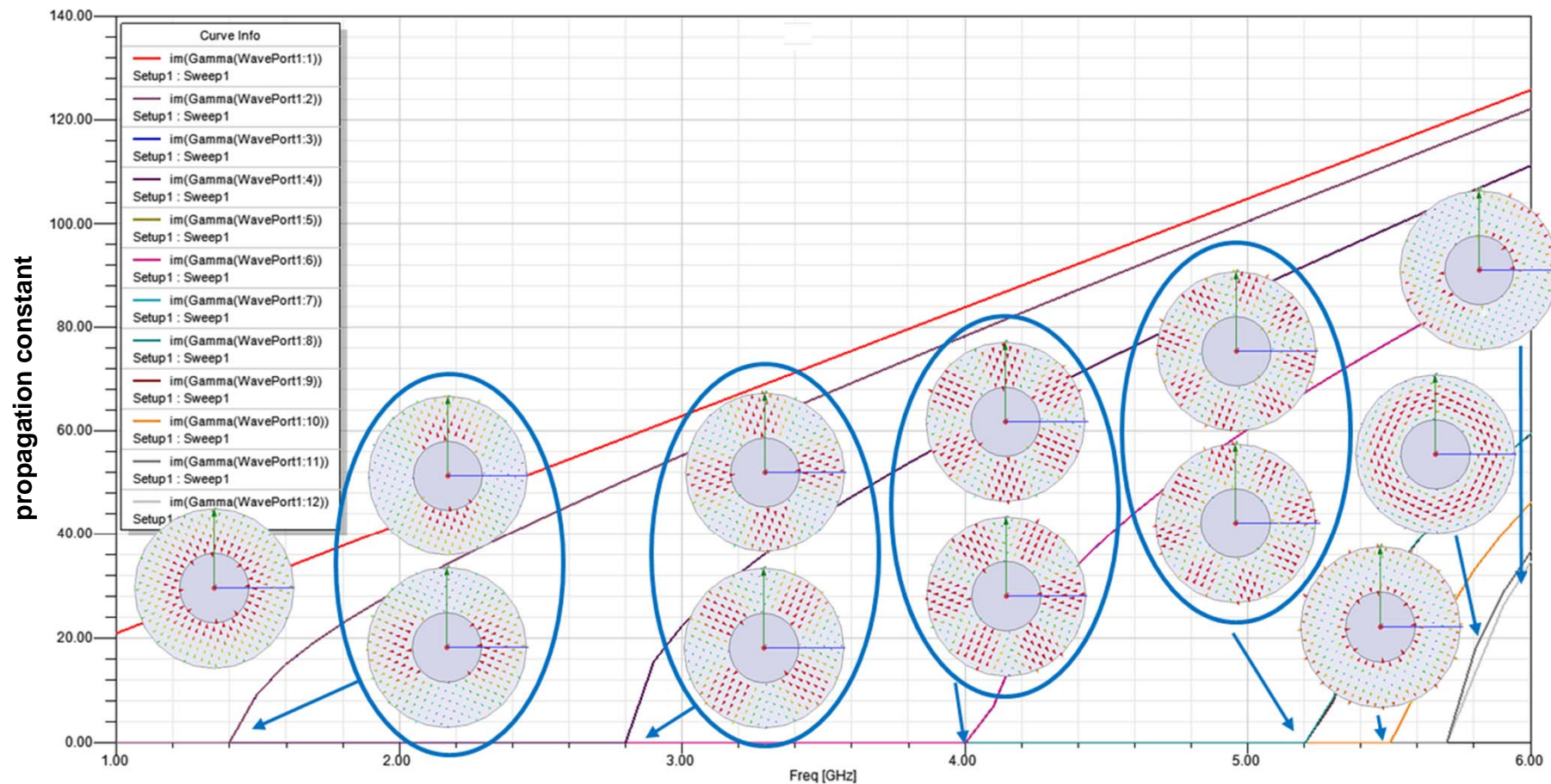
[Image from: <http://www.techlearner.com/Apps/TransandGuides.pdf>]

- The equations that rule the propagation of electromagnetic waves (Maxwell's equations) and the geometry of a given structure bound the electromagnetic field configurations that supports.
- These configurations are called (field) **modes**.

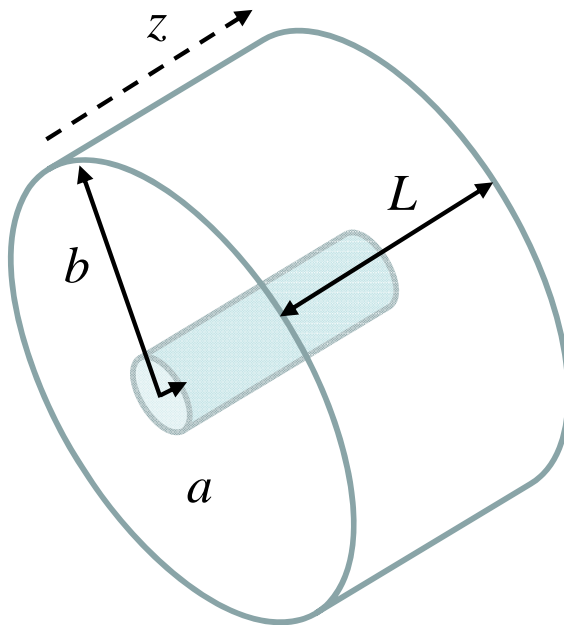




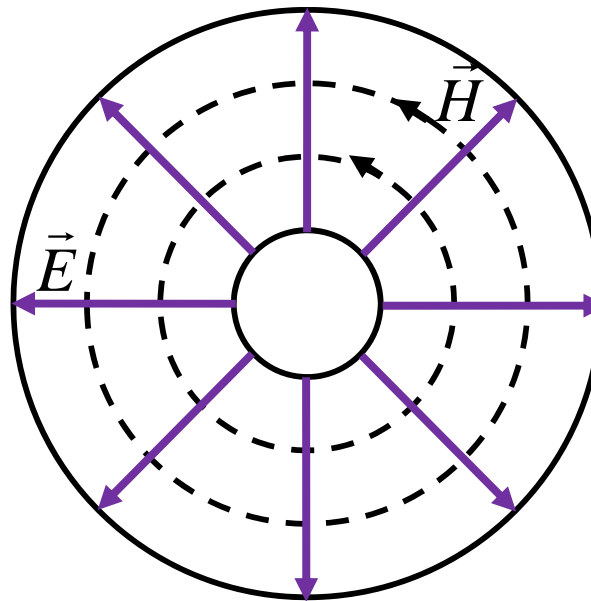
- The boundary conditions posed by the structures and set-ups (materials, voltage and/or current sources, frequencies) make possible that these modes do propagate or not inside the structures.



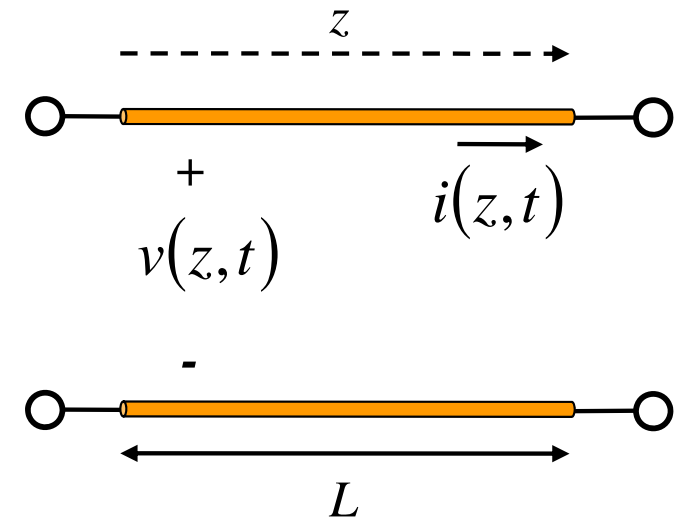
- A coaxial transmission line supports a TEM mode (electric field orientation, magnetic field orientation, and energy propagation direction form a triad).



physical structure



electromagnetic field distribution  
(TEM mode)

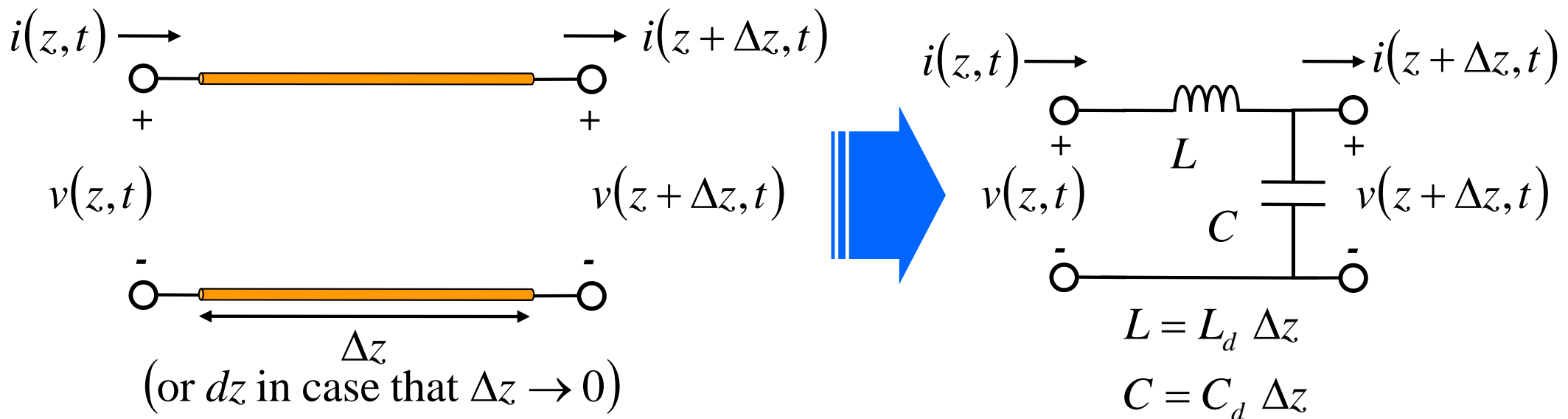


abstract model  
(ideal transmission line)

$$v(z, t) = - \int \vec{E}(z, t) d\vec{l}$$

$$i(z, t) = \oint \vec{H}(z, t) d\vec{l}$$

- The knowledge of voltage and current waves propagating along the transmission line allows the use of a distributed circuit model to analyze its performance.
- The model represents an infinitesimally short segment of the transmission line
- This model is convenient to explore properties of lines without knowing the fields in detail. However, the structures should be analyzed in detail if accurate performances have to be known.

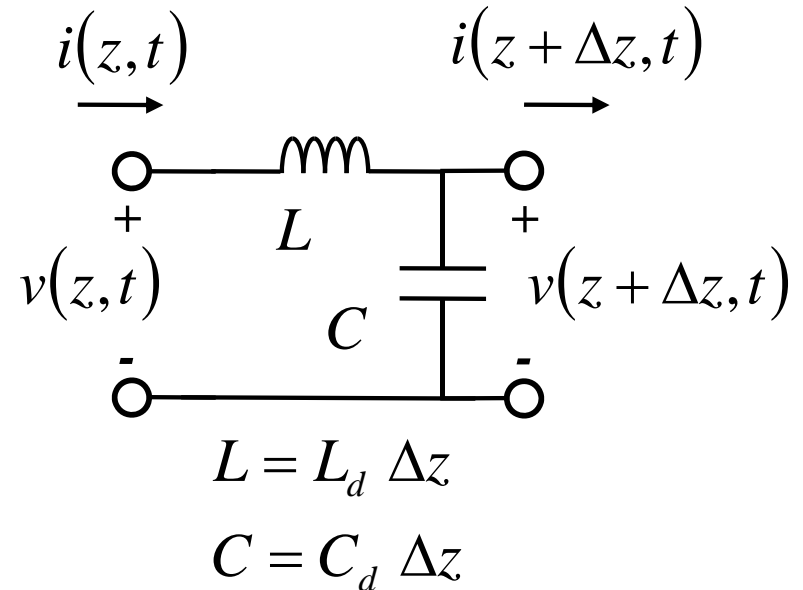


- $L_d$  and  $C_d$  are the distributed inductance [H/m] and capacitance [F/m] associated to the coaxial structure and materials. **No losses are assumed** in this example (meaning that there is no distributed resistance).

- Applying Kirchhoff's voltage and current laws:

$$v(z, t) - L_d \Delta z \frac{\partial i(z, t)}{\partial t} = v(z + \Delta z, t)$$

$$i(z, t) - C_d \Delta z \frac{\partial v(z + \Delta z, t)}{\partial t} = i(z + \Delta z, t)$$



- Dividing by  $\Delta z$  and taking the limit  $\Delta z \rightarrow 0$ :

$$\frac{\partial v(z, t)}{\partial z} = -L_d \frac{\partial i(z, t)}{\partial t}$$

$$\frac{\partial i(z, t)}{\partial z} = -C_d \frac{\partial v(z, t)}{\partial t}$$

- Considering sinusoidal **steady-state condition** (cosine based phasors) (**TRANSIENTS NOT CONSIDERED**):

$$\frac{\partial V}{\partial z} = -j\omega L_d I$$

$$\frac{\partial I}{\partial z} = -j\omega C_d V$$

- The **wave equations** can be solved simultaneously:

$$\frac{\partial^2 V}{\partial z^2} = -\beta^2 V$$

$$\frac{\partial^2 I}{\partial z^2} = -\beta^2 I$$

- Solutions are:

$$V(z) = V_0^+ e^{-j\beta z} + V_0^- e^{j\beta z}$$

$$I(z) = I_0^+ e^{-j\beta z} + I_0^- e^{j\beta z}$$

- Meaning that there are waves propagating in opposite directions along the transmission line (positive and negative waves).

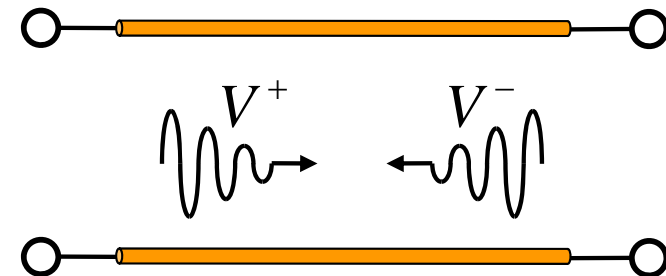
being:  $\beta = \omega \sqrt{L_d C_d}$

**phase constant**

being:  $I_0^+ = \frac{V_0^+}{Z_0}$        $I_0^- = -\frac{V_0^-}{Z_0}$

and:  $Z_0 = \sqrt{\frac{L_d}{C_d}}$

**characteristic impedance**



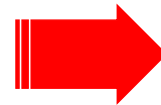


- Wave propagation solutions in time domain are:

$$v(z, t) = |V_0^+| \cos(\omega t - \beta z + \arg[V_0^+]) + |V_0^-| \cos(\omega t + \beta z + \arg[V_0^-])$$

$$i(z, t) = |I_0^+| \cos(\omega t - \beta z + \arg[I_0^+]) + |I_0^-| \cos(\omega t + \beta z + \arg[I_0^-])$$

the wave propagation on the line  
means a delay (in fact a phase delay)



lengths  $l$  are given in terms  
of  $\lambda$  or in degrees ( $\beta l$ )

- Wavelength and phase velocity on the line are:

$$\lambda = \frac{2\pi}{\beta} \quad v_p = \frac{\omega}{\beta} = \lambda f = \frac{1}{\sqrt{L_d C_d}}$$

$v_p$  does not change with  
frequency: NO DISTORTION  
(each frequency component of  
a signal travels at the same  $v_p$   
along the line)

- Do not forget... that the electrical model parameters **depend** on line **geometry**.



# Take your time...

**Characteristic impedance.** Given the solutions of the voltage and current wave equations derived for a two-wire lossless transmission line differential equations:

$$\begin{aligned}\frac{\partial V}{\partial z} &= -j\omega L_d I \\ \frac{\partial I}{\partial z} &= -j\omega C_d V\end{aligned} \quad \Rightarrow \quad \begin{aligned}V(z) &= V_0^+ e^{-j\beta z} + V_0^- e^{j\beta z} \\ I(z) &= I_0^+ e^{-j\beta z} + I_0^- e^{j\beta z}\end{aligned}$$

being:  $\beta = \omega \sqrt{L_d C_d}$

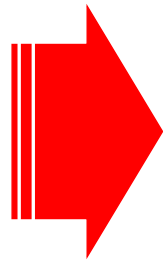
Find the characteristic impedance of the line and the relation between the voltage and current waves.

# Take your time...

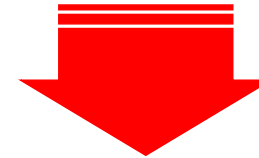
## Solution: Characteristic impedance.

Substituting the solution for the current into the differential equation for the voltage:

$$\frac{\partial V}{\partial z} = -j\omega L_d I$$



$$-j\beta V_0^+ e^{-j\beta z} + j\beta V_0^- e^{j\beta z} = -j\omega L_d (I_0^+ e^{-j\beta z} + I_0^- e^{j\beta z})$$



$$V(z) = V_0^+ e^{-j\beta z} + V_0^- e^{j\beta z}$$

$$I(z) = I_0^+ e^{-j\beta z} + I_0^- e^{j\beta z}$$

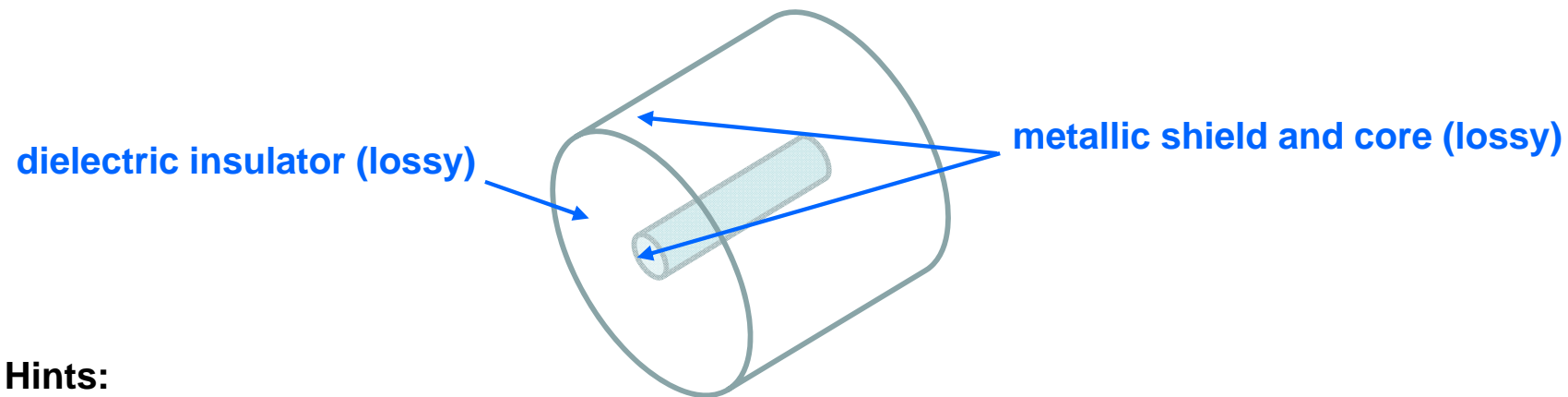
$$I_0^+ = \frac{V_0^+}{\omega L_d} \beta = \frac{V_0^+}{\sqrt{\frac{L_d}{C_d}}} = \frac{V_0^+}{Z_0}$$

characteristic  
impedance

$$I_0^- = -\frac{V_0^-}{Z_0}$$

# Take your time...

**Lossy two-wire transmission line.** In previous slides we presented the model of a two-wire lossless transmission line. Suggest a model for a lossy two-wire transmission line (e.g. lossy coaxial cable).



## Hints:

Losses in metals are characterized by their **electrical conductivity**  $\sigma$  (units:  $\Omega^{-1}\text{m}=\text{Sm}$ ) or by their **surface resistivity**  $R_s$  (units:  $\Omega/\text{square}$ ).  $R_s$  represents the mean power absorbed by a unit area ( $1 \text{ m}^2$ ) and can be used on planar surfaces or surfaces with curvature radius smaller than the **skin depth**  $\delta_s$ . Currents on metals flow mainly in their outer “skin” at a level called skin depth.

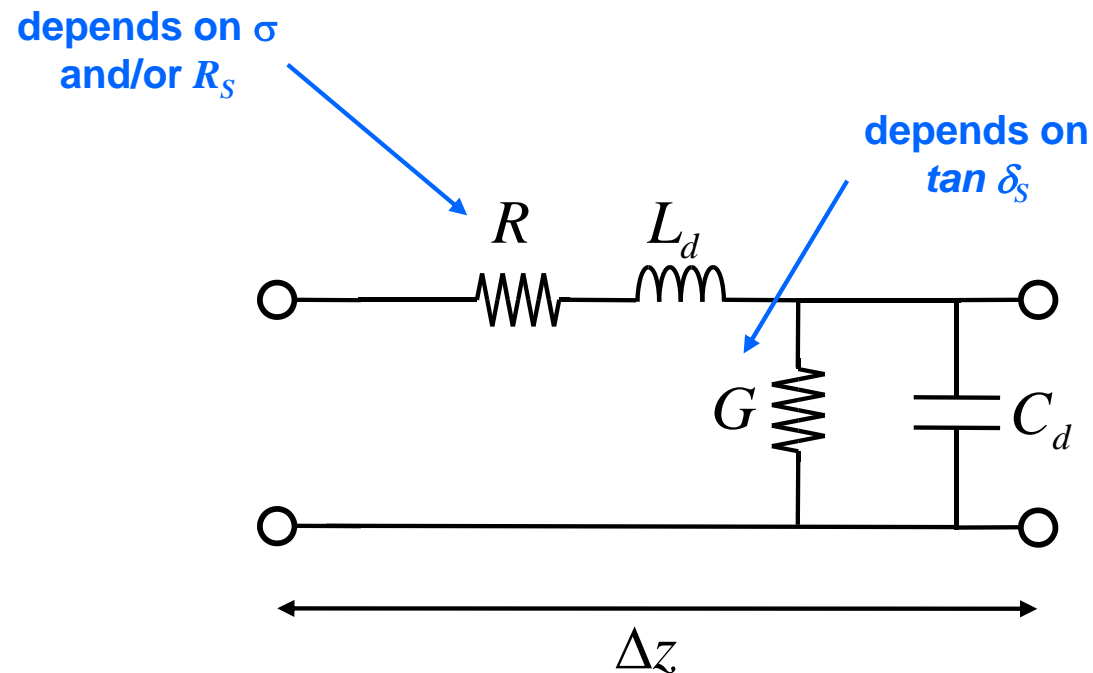
Losses in dielectrics are characterized by their **loss tangent**  $\tan \delta$ . It represents the tangent of the angle in the complex plane between the electric field resistive losses and its lossless reactive component.

$$R_s = \frac{1}{\delta_s \sigma}$$
$$\delta_s = \frac{1}{\sqrt{\pi \mu \sigma f}}$$

$$\tan \delta = \frac{\omega \epsilon'' + \sigma_d}{\omega \epsilon'}$$

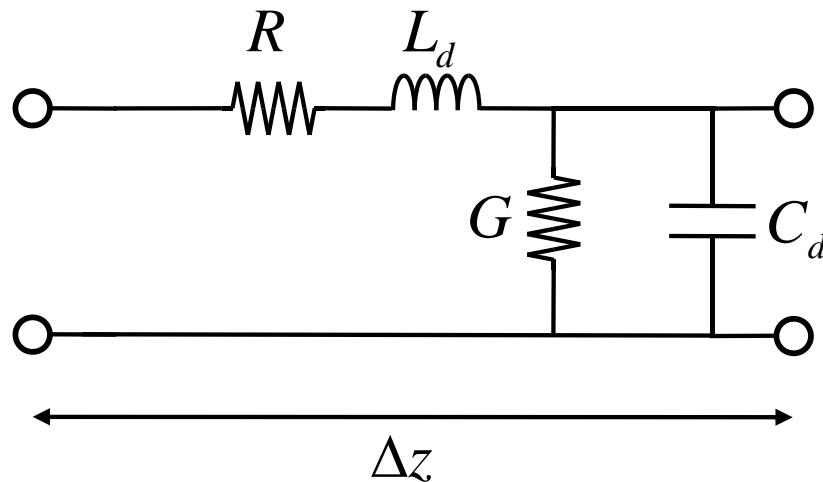
# Take your time...

Solution. Lossy two-wire transmission line.



# Take your time...

**Telegrapher's equations.** Derive the propagation equations for a lossy two-wire transmission line having the lumped-circuit model of the figure. The resulting set of differential equations are called the **telegrapher's equations** and are due **Oliver Heaviside** in 1880.

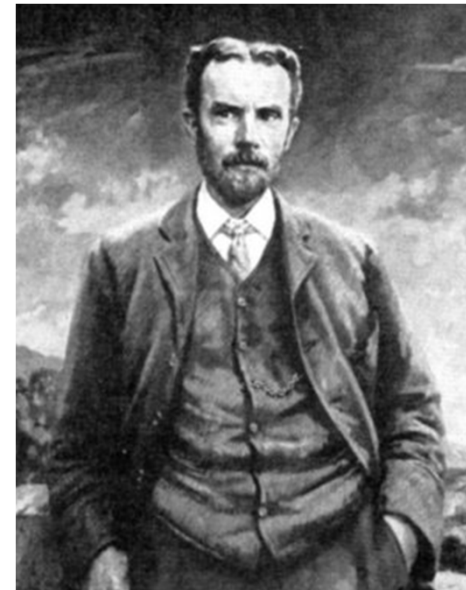


$L_d$  distributed inductance per unit length [H/m]

$C_d$  distributed capacitance per unit length [F/m]

$R$  distributed resistance per unit length, for both conductors [ $\Omega$ /m]

$G$  distributed conductance per unit length, for both conductors [S/m]



[Image from:  
[http://es.wikipedia.org/wiki/Oliver\\_Heaviside](http://es.wikipedia.org/wiki/Oliver_Heaviside)]

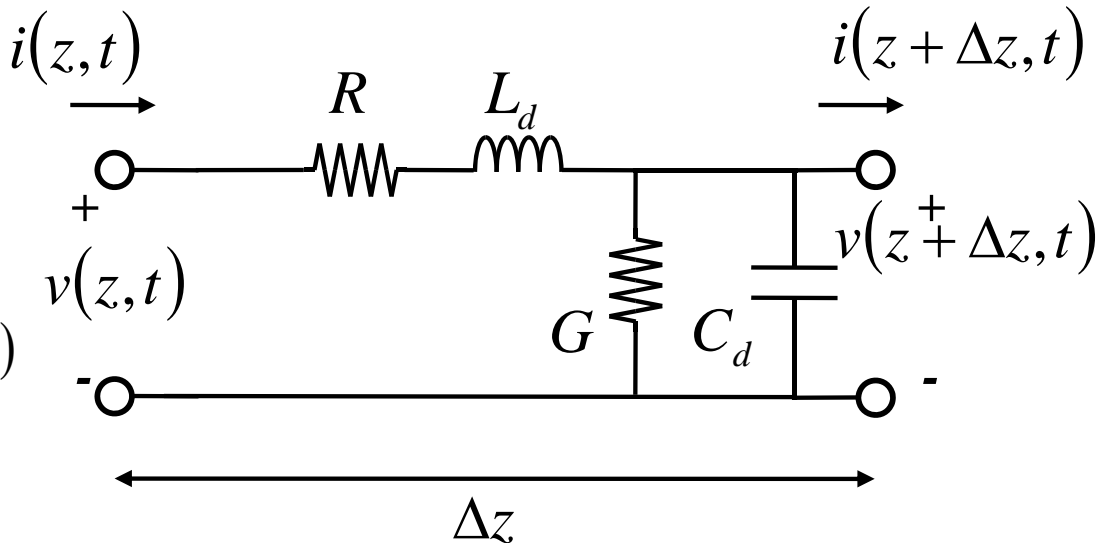
# Take your time...

## Solution: Telegrapher's equations.

- According to Kirchhoff's laws:

$$v(z,t) - L_d \Delta z \frac{\partial i(z,t)}{\partial t} - R \Delta z i(z,t) = v(z + \Delta z, t)$$

$$i(z,t) - G \Delta z v(z + \Delta z, t) - C_d \Delta z \frac{\partial v(z + \Delta z, t)}{\partial t} = i(z + \Delta z, t)$$



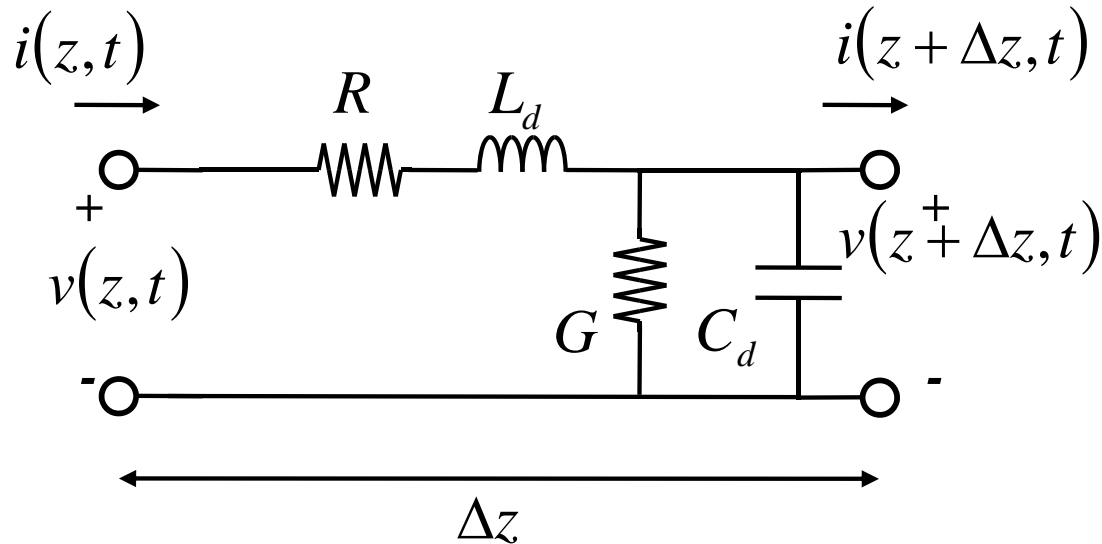
- Dividing by  $\Delta z$  and taking the limit  $\Delta z \rightarrow 0$ :

$$\frac{\partial v(z,t)}{\partial z} = -R i(z,t) - L_d \frac{\partial i(z,t)}{\partial t}$$

$$\frac{\partial i(z,t)}{\partial z} = -G v(z,t) - C_d \frac{\partial v(z,t)}{\partial t}$$

## Telegrapher's equations

- From the lumped-circuit model of a **lossy** two-wire transmission line we get the telegrapher's equations.



$$\frac{\partial v(z, t)}{\partial z} = -R i(z, t) - L_d \frac{\partial i(z, t)}{\partial t}$$

$$\frac{\partial i(z, t)}{\partial z} = -G v(z, t) - C_d \frac{\partial v(z, t)}{\partial t}$$

- Considering sinusoidal steady-state condition (cosine based phasors):

$$\frac{\partial V}{\partial z} = -(R + j\omega L_d) I$$

$$\frac{\partial I}{\partial z} = -(G + j\omega C_d) V$$



- Wave equations are now:

$$\frac{\partial^2 V}{\partial z^2} = \gamma^2 V$$

$$\frac{\partial^2 I}{\partial z^2} = \gamma^2 I$$

being:  $\gamma = \alpha + j\beta = \sqrt{(R + j\omega L_d)(G + j\omega C_d)}$

complex propagation constant    attenuation    phase constant

$v_p$  depends with frequency: **DISTORTION**  
(each frequency component of a signal travels at different  $v_p$  along the line)

- The solutions are:

$$V(z) = V_0^+ e^{-\gamma z} + V_0^- e^{\gamma z}$$

$$I(z) = I_0^+ e^{-\gamma z} + I_0^- e^{\gamma z} = \frac{1}{Z_0} (V_0^+ e^{-\gamma z} - V_0^- e^{\gamma z})$$

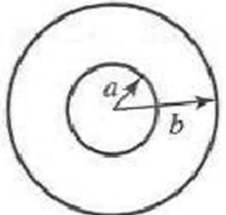
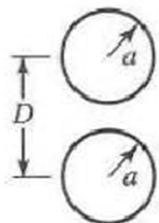
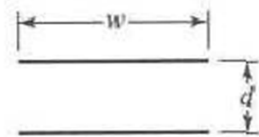
being:  $Z_0 = \sqrt{\frac{R + j\omega L_d}{G + j\omega C_d}}$

- And in time domain:

$$v(z, t) = |V_0^+| e^{-\alpha z} \cos(\omega t - \beta z + \arg[V_0^+]) + |V_0^-| e^{\alpha z} \cos(\omega t + \beta z + \arg[V_0^-])$$

- Lumped-circuit parameters required to model some common lines as a function of their dimensions, surface resistivity ( $R_s$ ), and materials filling the space between the conductors (permittivity  $\epsilon = \epsilon' - j\epsilon''$  and permeability  $\mu = \mu_0\mu_r$ ).

TABLE 2.1 Transmission Line Parameters for Some Common Lines

	COAX	TWO-WIRE	PARALLEL PLATE
			
$L$	$\frac{\mu}{2\pi} \ln \frac{b}{a}$	$\frac{\mu}{\pi} \cosh^{-1} \left( \frac{D}{2a} \right)$	$\frac{\mu d}{w}$
$C$	$\frac{2\pi\epsilon'}{\ln b/a}$	$\frac{\pi\epsilon'}{\cosh^{-1}(D/2a)}$	$\frac{\epsilon' w}{d}$
$R$	$\frac{R_s}{2\pi} \left( \frac{1}{a} + \frac{1}{b} \right)$	$\frac{R_s}{\pi a}$	$\frac{2R_s}{w}$
$G$	$\frac{2\pi\omega\epsilon''}{\ln b/a}$	$\frac{\pi\omega\epsilon''}{\cosh^{-1}(D/2a)}$	$\frac{\omega\epsilon'' w}{d}$

parameters for other transmission lines are found in the literature or by using specific software

[Table from: [POZAR]]

- In practical lines losses are small. The equations for attenuation and propagation factor can be simplified.

- The complex propagation factor can be re-arranged:

$$\gamma = \alpha + j\beta = \sqrt{(R + j\omega L_d)(G + j\omega C_d)} = j\omega\sqrt{L_d C_d} \sqrt{1 - j\left(\frac{R}{\omega L_d} + \frac{G}{\omega C_d}\right) - \frac{RG}{\omega^2 L_d C_d}}$$

- For a low-loss line we can assume:  $R \ll \omega L_d$  and  $G \ll \omega C_d$

$$\gamma \approx j\omega\sqrt{L_d C_d} \sqrt{1 - j\left(\frac{R}{\omega L_d} + \frac{G}{\omega C_d}\right)} \approx j\omega\sqrt{L_d C_d} \left[1 - \frac{j}{2}\left(\frac{R}{\omega L_d} + \frac{G}{\omega C_d}\right)\right]$$

$$Z_0 = \sqrt{\frac{R + j\omega L_d}{G + j\omega C_d}} \approx \sqrt{\frac{L_d}{C_d}}$$

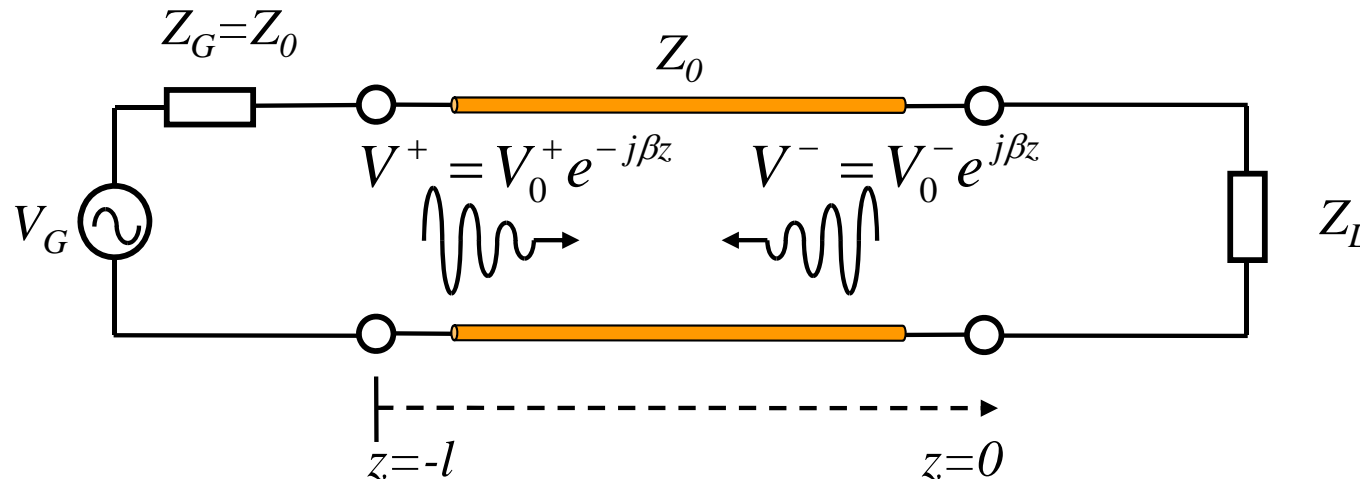
- Finally:  $\alpha \approx \frac{1}{2}\left(R\sqrt{\frac{C_d}{L_d}} + G\sqrt{\frac{L_d}{C_d}}\right)$

$$\beta \approx \omega\sqrt{L_d C_d} \quad Z_0 \approx \sqrt{\frac{L_d}{C_d}}$$

same values than lossless lines

**DISTORTIONLESS**

- Wave reflection on a transmission line can be illustrated by considering a lossless transmission line loaded with an arbitrary impedance  $Z_L$ .  $Z_0$  is the characteristic impedance of the transmission line.



- A voltage reflection coefficient can be defined for any point in the line as the amplitude of the reflected voltage wave normalized to the amplitude of the incident voltage wave.
- Because **at the load** ( $z=0$ ) the impedance of the line is  $Z_L$ :

$$\rho(z) = \frac{V^-(z)}{V^+(z)}$$

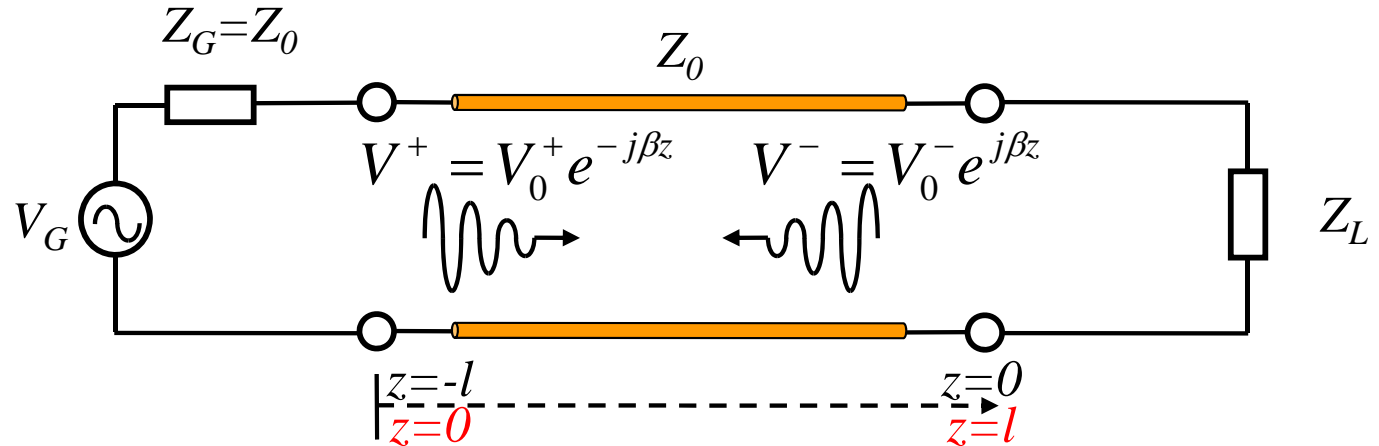
$$Z_L = \frac{V(0)}{I(0)} = \frac{V_0^+ + V_0^-}{V_0^+ - V_0^-} Z_0 \quad \Rightarrow \quad \rho_L = \rho(0) = \frac{Z_L - Z_0}{Z_L + Z_0}$$

**voltage reflection coefficient at load**

- One has to be careful with the coordinate axis chosen to define the reflection coefficient along the line.

- by definition:

$$\rho(z) = \frac{V^-(z)}{V^+(z)}$$



- coefficient at load:  $\rho_L = \rho(0) = \frac{V_0^-}{V_0^+} = \frac{Z_L - Z_0}{Z_L + Z_0}$   $\rho_L = \rho(l) = \frac{V_0^-}{V_0^+} e^{j2\beta l} = \frac{Z_L - Z_0}{Z_L + Z_0}$

- coefficient at input port:  $\rho_{IN} = \rho(-l) = \frac{V_0^-}{V_0^+} e^{-j2\beta l}$   $\rho_{IN} = \rho(0) = \frac{V_0^-}{V_0^+}$

- in both cases:

$$\rho_{IN} = \rho_L e^{-j2\beta l}$$

- however:

$$\rho(z) = \rho_L e^{j2\beta z}$$

$$\rho(z) = \rho_L e^{-j2\beta(l-z)}$$

- Because **at the input port of the line** ( $z=-l$ ) the impedance is:

$$\rho_{IN} = \rho(-l) = \frac{V^-(-l)}{V^+(-l)} = \frac{V_0^- e^{-j\beta l}}{V_0^+ e^{j\beta l}} = \rho(0) e^{-2j\beta l} = \rho_L e^{-2j\beta l}$$

$$\rho_{IN} = \rho_L e^{-2j\beta l}$$

**voltage reflection coefficient at input port**

- At any point in the line:  $\rho(z) = \rho_L e^{2j\beta z}$
- Reflection coefficient is a complex number.
- For a passive load the magnitude of the reflection coefficient is always lower than 1.

$$0 \leq |\rho| \leq 1$$

**Example: Reflection coefficient of a loaded **lossy** transmission line.** Taking into consideration the equations of the voltage and current waves flowing in a lossy transmission line find the equation for the input impedance on the line.

$$\rho_{IN} = \rho(-l) = \frac{V^-(-l)}{V^+(-l)} = \frac{V_0^- e^{-\gamma l}}{V_0^+ e^{\gamma l}} = \rho_L e^{-2\alpha l} e^{-2j\beta l}$$

The reflection coefficient is attenuated when the line increases its length.

# Take your time...

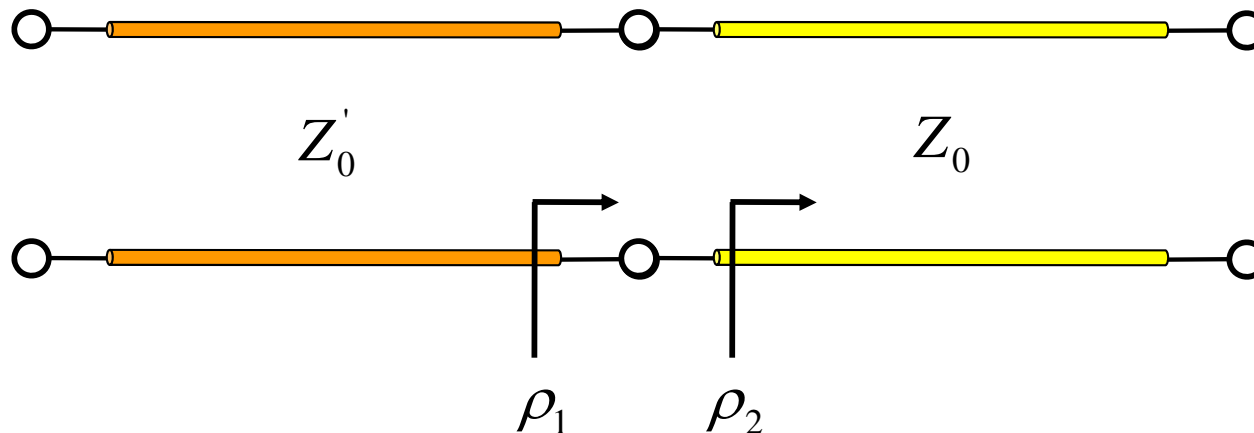
**Standard loads.** Which is the reflection coefficient corresponding to, respectively, an open circuit, short circuit, and reference impedance ( $Z_0$ )?

$$\rho_L|_{Z_L \rightarrow \infty}$$

$$\rho_L|_{Z_L=0}$$

$$\rho_L|_{Z_L=Z_0}$$

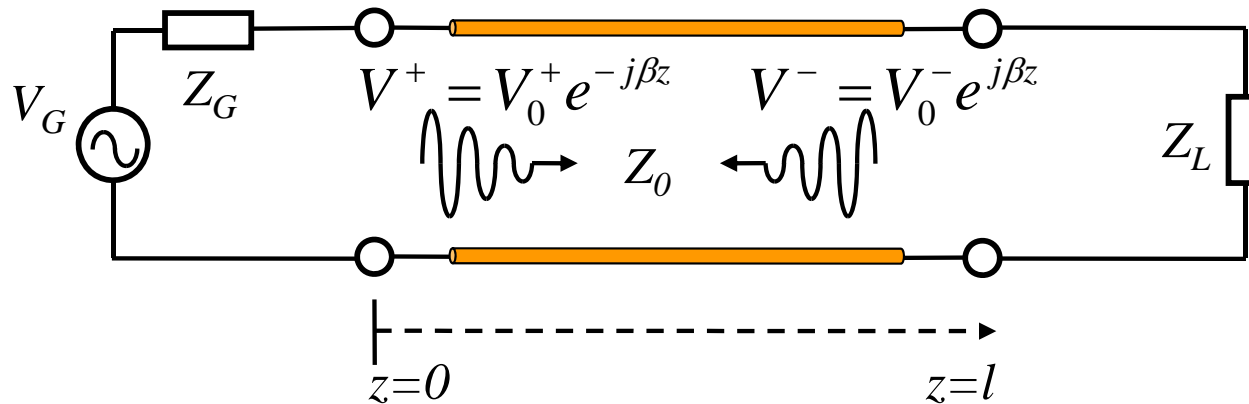
**Line transition.** Is the voltage reflection coefficient the same at both sides of the transitions between the transmission lines?





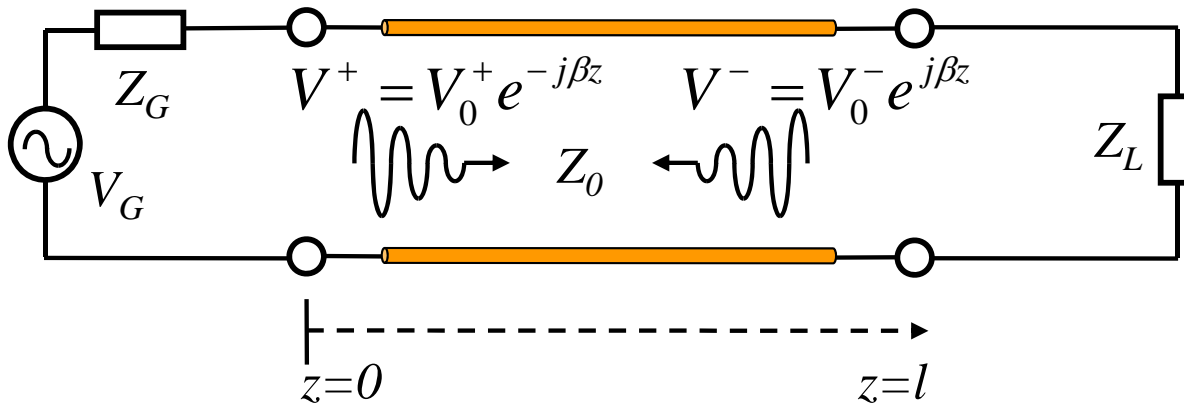
# Take your time...

**Mismatched lossless transmission line.** By applying boundary conditions to the ports of the lossless transmission line of the figure, find the voltage reflection coefficients  $\rho_{IN}$  and  $\rho_L$ , and the magnitude of the progressive wave  $V_0^+$ .



# Take your time...

## Solution. Mismatched lossless transmission line.



- at the output port of the line:

- voltage:

$$V^+(l) + V^-(l) = V_0^+ e^{-j\beta l} + V_0^- e^{j\beta l} = I(0)Z_L = \frac{Z_L}{Z_0} [V_0^+ e^{-j\beta l} - V_0^- e^{j\beta l}]$$

- defining the load reflection coefficient  $\rho_L$ ...

$$\rho_L = \frac{V^-(l)}{V^+(l)} = \frac{V_0^-}{V_0^+} e^{2j\beta l}$$

- ... we get:

$$\rho_L = \frac{Z_L - Z_0}{Z_L + Z_0}$$

- and at the input port:

- voltage and current:

$$\left. \begin{aligned} V^+(0) + V^-(0) &= V_G - I(0)Z_G \\ \frac{V^+(0)}{Z_0} - \frac{V^-(0)}{Z_0} &= I(0) \end{aligned} \right\}$$

- defining  $\rho_{IN}$ ...

$$\rho_{IN} = \frac{V^-(0)}{V^+(0)} = \rho_L e^{-j2\beta l}$$

- and  $\rho_G$ ...

$$\rho_G = \frac{Z_G - Z_0}{Z_G + Z_0}$$

- ...we get:

$$V_0^+ = V_G \frac{Z_0}{Z_0 + Z_G} \frac{1}{1 - \rho_{IN}\rho_G}$$

- The time-average power flow along the line at point  $z$  is:

$$P_{av} = \frac{1}{2} \operatorname{Re}[V(z)I^*(z)] = \frac{|V_0^+|^2}{2Z_0} - \frac{|V_0^-|^2}{2Z_0} = P^+ - P^-$$

$$P_{av} = \frac{|V_0^+|^2}{2Z_0} (1 - |\rho(z)|^2)$$

- Incident power to the line:

$$P_{inc} = \frac{|V_0^+|^2}{2Z_0}$$

- Reflected power at the load:

$$P_{ref} = \frac{|V_0^+|^2}{2Z_0} |\rho_L|^2$$

- To avoid the existence of reflected waves ( $\rho=0$ ) on a loaded transmission line the load impedance  $Z_L$  should be equal to the characteristic impedance  $Z_0$ . Such a load is said to be a **matched load**.

- Some considerations derive from the previous equations: power on the load

$$\begin{array}{ccc} Z_L = Z_0 & \longrightarrow & P^- = 0 \quad P_L = P^+ \\ Z_L = \{0, \infty\} & & P^- = P^+ \quad P_L = 0 \end{array}$$

- The power carried by the positive flowing wave can be greater than the average power flowing on the line
- When the load is mismatched not all the available power is delivered to the load. This loss is called **return loss** (RL):

$$RL = -20 \cdot \log |\rho_L| \quad [\text{dB}]$$

- What happens in case of a **lossy** transmission line?

For a forward propagating wave in a transmission line with length  $l$ :

$$\begin{aligned}
 V(z) &= V_0^+ e^{-\gamma z} \\
 I(z) &= \frac{1}{Z_0} V_0^+ e^{-\gamma z}
 \end{aligned}
 \quad \Rightarrow \quad
 P^+(l) = \frac{1}{2} \operatorname{Re}[V \cdot I^*] = P^+(0) e^{-2\alpha l}$$

attenuation  
 incoming power at input port

Wave attenuation ( $L$ ) between planes  $z=0$  and  $z=l$ :

- in decibels [dB]: 
$$L [\text{dB}] = 10 \log \left( \frac{P^+(0)}{P^+(l)} \right) = \alpha l \, 20 \log e = 8.686 \, \alpha l$$

# Take your time...

**Feeding and antenna.** A  $50\text{-}\Omega$  antenna is fed with a signal of  $100\text{ W}$  by means of a  $75\text{ }\Omega$  cable. Find the power finally arriving to the antenna.

- a)  $4\text{ W}$       b)  $20\text{ W}$       c)  $80\text{ W}$       d)  $96\text{ W}$

- The voltage on each point of a line depends on the load attached at its end:

$$|V(z)| = |V_0^+| |1 + \rho_L e^{j2\beta z}|$$

- When the transmission line is matched, the magnitude of the voltage on the line is constant:

$$|V(z)| = |V_0^+|$$

- When the transmission line is not matched the overlap of an incoming and a reflected wave leads to a **standing wave** whose magnitude oscillates with the position on the line:

$$V_{\max} = |V_0^+| (1 + |\rho_L|)$$

$$V_{\min} = |V_0^+| (1 - |\rho_L|)$$

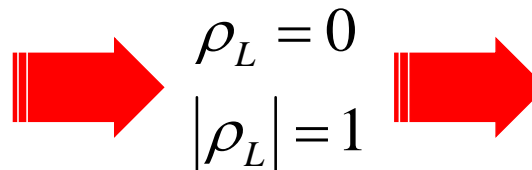
- A measure of the line mismatch is the **Voltage Standing Wave Ratio** (VSWR):

$$VSWR = \frac{V_{\max}}{V_{\min}} = \frac{1 + |\rho_L|}{1 - |\rho_L|}$$

- SWR is a real positive number:  $1 \leq VSWR \leq \infty$

matched load

total reflection

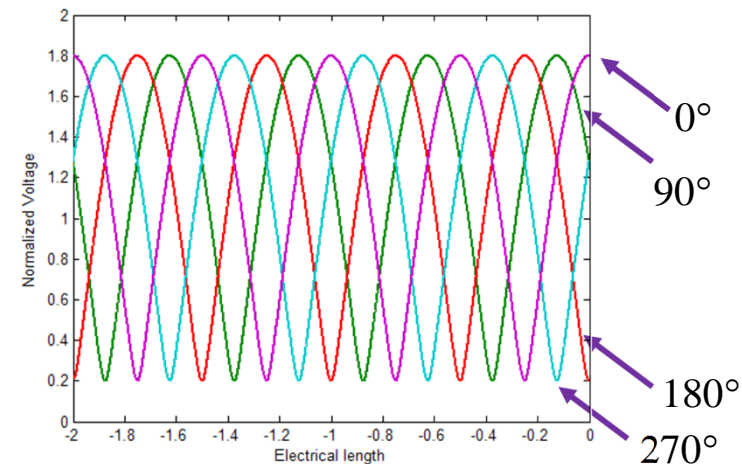
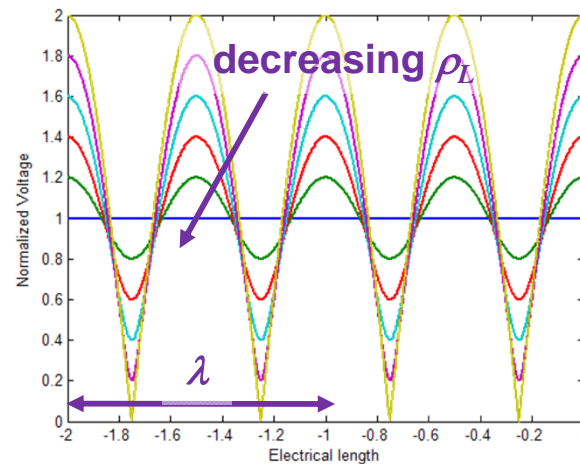


$$VSWR = 1$$

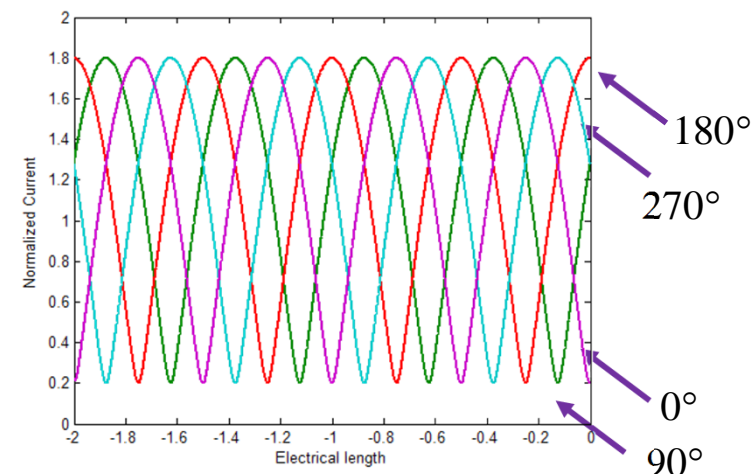
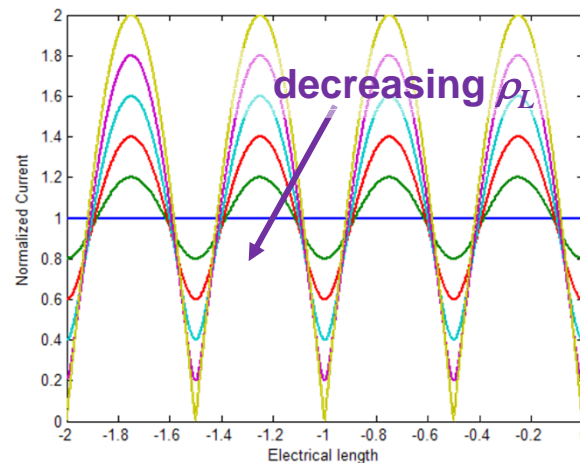
$$VSWR = \infty$$

**Example: Voltage on a transmission line.** Plot the normalized voltage and current on a transmission line  $2\lambda$  in length as a function of position on the line ( $z$ ). Consider the following loads:  $\rho_L = \{0, 0.2, 0.4, 0.6, 0.8, 1\}$  and  $\rho_L = \{0.8, 0.8 e^{j90^\circ}, 0.8 e^{j180^\circ}, 0.8 e^{j270^\circ}\}$ .

$$\frac{|V(z)|}{|V_0^+|} = |1 + \rho_L e^{j2\beta z}|$$



$$\frac{|I(z)|}{|I_0^+|} = |1 - \rho_L e^{j2\beta z}|$$



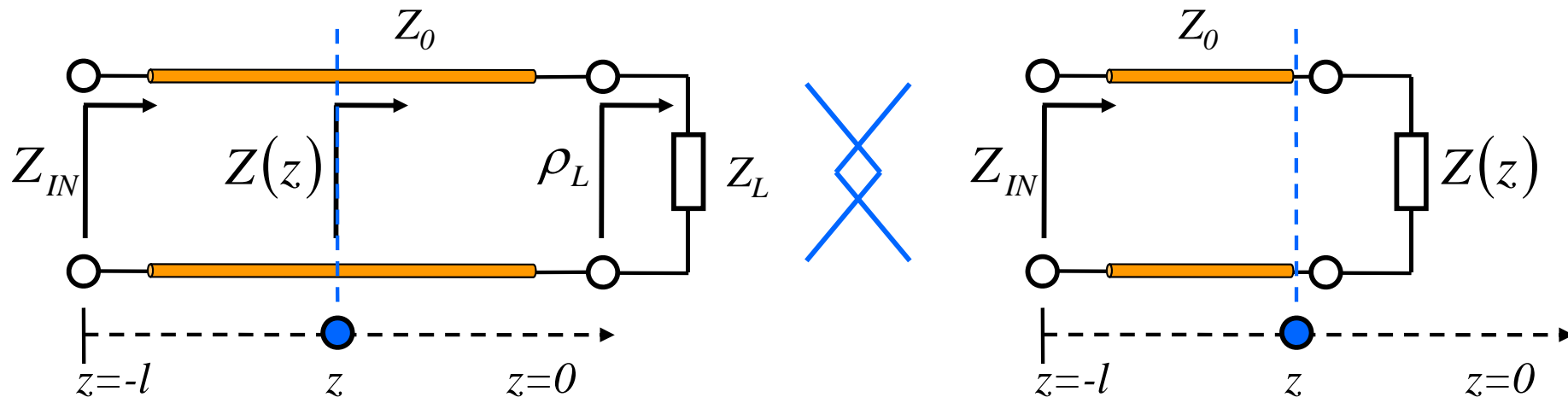


- Line impedance varies with position ( $z$ ):


$$Z(z) = \frac{V(z)}{I(z)} = Z_0 \frac{V_0^+ e^{-j\beta z} + V_0^- e^{j\beta z}}{V_0^+ e^{-j\beta z} - V_0^- e^{j\beta z}} = Z_0 \frac{1 + \rho_L e^{j2\beta z}}{1 - \rho_L e^{j2\beta z}}$$


- The input impedance of a line loaded with  $Z_L$  is:

$$Z_{IN} = \frac{V(-l)}{I(-l)} = Z_0 \frac{Z_L \cos(\beta l) + jZ_0 \sin(\beta l)}{Z_0 \cos(\beta l) + jZ_L \sin(\beta l)} = Z_0 \frac{Z_L + jZ_0 \tan(\beta l)}{Z_0 + jZ_L \tan(\beta l)}$$




- The case of highly reflective loads:


short-circuited stub:  $Z_{IN} = jZ_0 \tan(\beta l)$   inductive impedance

open-circuited stub:  $Z_{IN} = \frac{Z_0}{j \tan(\beta l)}$   capacitive impedance

- Quarter-wave long transmission lines perform as **quarter-wave (impedance) transformers** or **impedance inverters**, being the input impedance inversely proportional to the load impedance.

When:  $l = \frac{\lambda}{4} + n \frac{\lambda}{2} \quad n = 0, 1, 2, \dots$    $Z_{IN} = \frac{Z_0^2}{Z_L}$

- Otherwise, any line whose length is any multiple of  $\lambda/2$  does not transform the load impedance, regardless of the characteristic impedance.

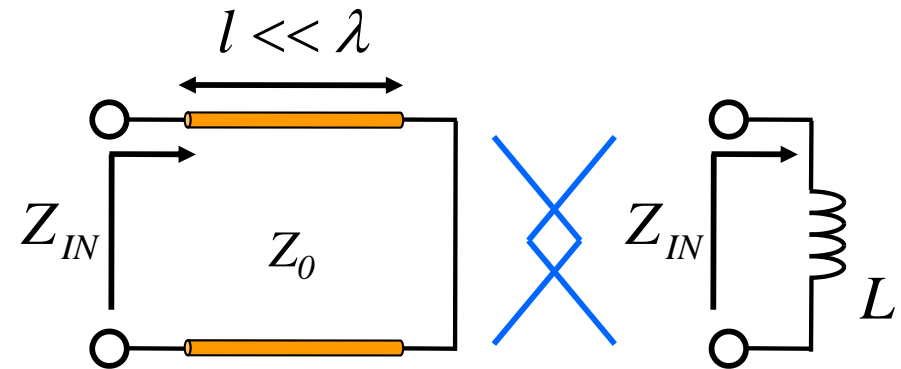
When:  $l = n \frac{\lambda}{2} \quad n = 1, 2, 3, \dots$    $Z_{IN} = Z_L$

### Example: Emulating lumped elements with electrically-short transmission lines.

Suggest what lumped elements can be simulated by means of short transmission lines when these line end, respectively, with an open circuit and with a short circuit.

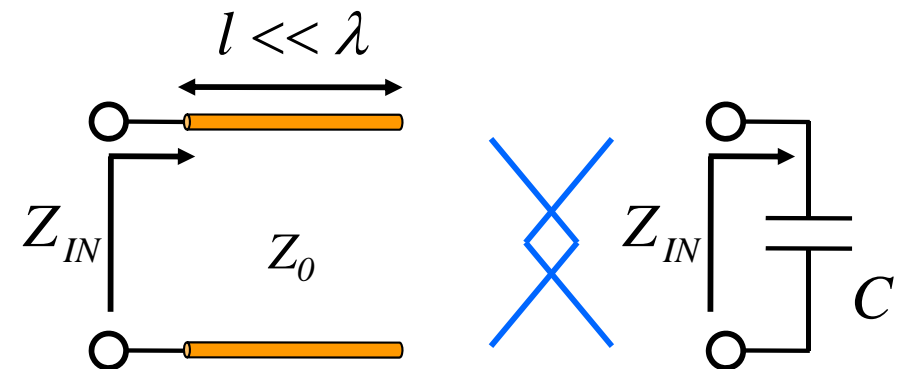
short-circuited short-stub:

$$Z_{IN} \approx jZ_0\beta l = j\omega L \quad \text{being: } L = \frac{Z_0 l}{v_p}$$

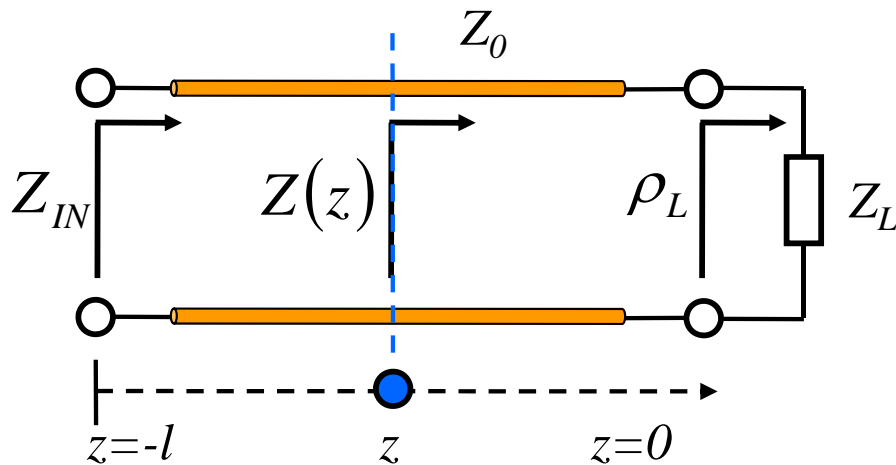


open-circuited short-stub:

$$Z_{IN} \approx \frac{Z_0}{j\beta l} = \frac{1}{j\omega C} \quad \text{being: } C = \frac{v_p l}{Z_0}$$



**Example: Input impedance of a loaded **lossy** transmission line.** Taking into consideration the equations of the voltage and current waves flowing in a lossy transmission line, find the equation for the input impedance on the line.



$$V(z) = V_0^+ e^{-\gamma z} + V_0^- e^{\gamma z}$$

$$I(z) = I_0^+ e^{-\gamma z} + I_0^- e^{\gamma z} = \frac{1}{Z_0} (V_0^+ e^{-\gamma z} - V_0^- e^{\gamma z})$$


 at  $z = -l$

that is, also

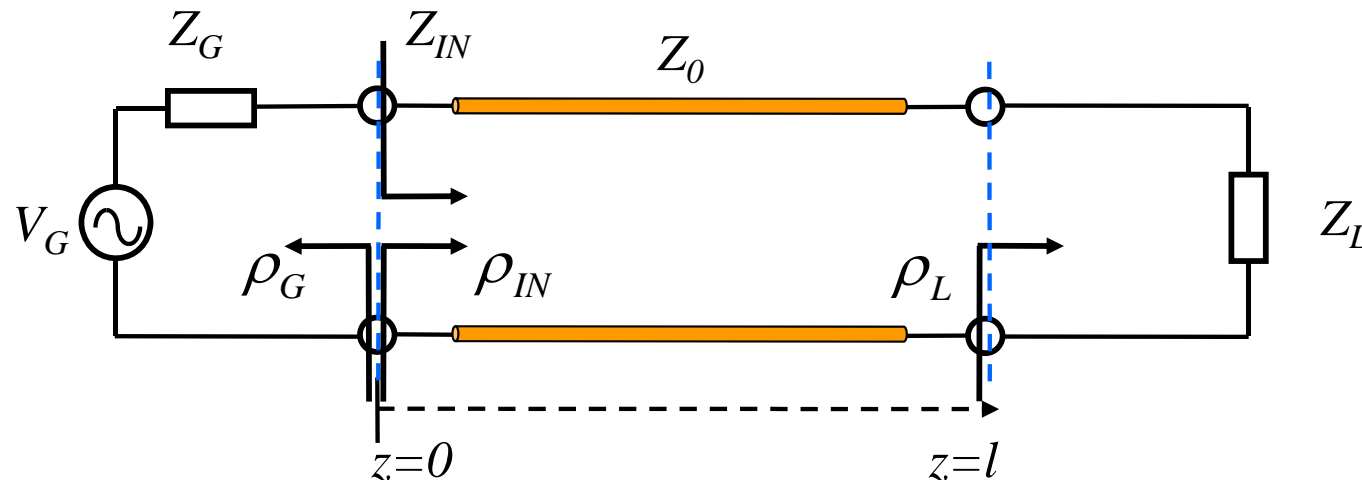
$$Z_{IN} = Z_0 \tanh(\gamma l + \delta_Z)$$

$$\delta_Z = \tanh^{-1} \left( \frac{Z_L}{Z_0} \right)$$



$$Z_{IN} = Z_0 \frac{Z_L \cosh(\gamma l) + Z_0 \sinh(\gamma l)}{Z_0 \cosh(\gamma l) + Z_L \sinh(\gamma l)}$$

- The voltage on the input port of the line can be calculated from the source voltage and depends on the load impedance.





$$V_0^+ = V_G \frac{Z_0}{Z_0 + Z_G} \frac{1}{1 - \rho_{IN} \rho_G} \quad \text{being:} \quad \rho_G = \frac{Z_G - Z_0}{Z_G + Z_0}$$

at  $z=0$ : 
$$V(0) = V_G \frac{Z_{IN}}{Z_{IN} + Z_G} = V_0^+ + V_0^- = V_0^+ (1 + \rho_{IN})$$

and being: 
$$Z_{IN} = \frac{V(0)}{I(0)} = Z_0 \frac{1 + \rho_{IN}}{1 - \rho_{IN}}$$

- Consequently:

when:  $Z_G = Z_0$    $V_0^+ = \frac{V_G}{2}$

$Z_{IN} = Z_0$    $V_0^+ = V_G \frac{Z_0}{Z_0 + Z_G}$

the power delivered from the  
source to the line (and if it has  
no losses, delivered to the load):

$$P_{av} = \frac{1}{2} \operatorname{Re}[V(0)I^*(0)] = \frac{|V_G|^2}{2} \left| \frac{Z_{IN}}{Z_{IN} + Z_G} \right|^2 \operatorname{Re}\{Z_{IN}\}$$

being:  $Z_{IN} = R_{IN} + jX_{IN}$

and:  $Z_G = R_G + jX_G$

$$P_{av} = \frac{|V_G|^2}{2} \frac{R_{IN}}{(R_{IN} + R_G)^2 + (X_{IN} + X_G)^2}$$

# Take your time...

**Power delivered to the load.** Consider a lossless transmission line connected to a fixed source impedance  $Z_G = R_G + jX_G$ . Find the power delivered to the load (or the power delivered to the transmission line) in the following two cases: a) when the load is matched to the line ( $Z_L = Z_0$ ); and b) when the generator is matched to the loaded line ( $Z_G = Z_{IN}$ ).

# Take your time...

**Solution: Power delivered to the load.**

- Load matched to the line ( $Z_L=Z_0$ ).

$$P_{av} = \frac{|V_G|^2}{2} \frac{Z_0}{(Z_0 + R_G)^2 + X_G^2}$$

- Generator matched to the loaded line ( $Z_G=Z_{IN}$ ).

$$P_{av} = \frac{|V_G|^2}{8} \frac{R_G}{R_G^2 + X_G^2}$$



# Take your time...

**Impedance for maximum power transfer or available power.** Assuming that the generator series impedance is fixed find the input impedance  $Z_{IN}$  to achieve the maximum power transfer to the load (lossless transmission line). In that case, find the power delivered to the load.

# Take your time...

**Solution: Impedance for maximum power transfer.**

- Maximizing the power transfer...

$$\left. \begin{array}{l} \frac{\partial P_{av}}{\partial R_{IN}} = 0 \\ \frac{\partial P_{av}}{\partial X_{IN}} = 0 \end{array} \right\} \Rightarrow \left. \begin{array}{l} R_G^2 - R_{IN}^2 + (X_G + X_{IN})^2 = 0 \\ R_{IN}(X_G + X_{IN}) = 0 \end{array} \right\} \Rightarrow \left. \begin{array}{l} R_{IN} = R_G \\ X_{IN} = -X_G \end{array} \right\} \Rightarrow Z_{IN} = Z_G^*$$

... we get that the input impedance of the line should be the complex conjugate of the source impedance. This condition is called **conjugate matching**.

- In this case, the power transferred to the load is (lossless transmission line):

$$P_{av} = \frac{|V_G|^2}{8R_G}$$